

BEARCLAW *Hermes* BREACH MODEL

PURPOSE	2D model of <i>Hermes</i> spaceship breach
EQUATION	Euler equations $q_t + f(q)_x + f(q)_y = 0$
DOMAIN	$-100 \leq x \leq 100, -20 \leq y \leq 20$
INITIAL CONDITION	$p = p_0$ for $x > 0$, $p = p_v \ll p_0$ for $x \leq 0$
BOUNDARY CONDITIONS	Characteristic outflow, moving wall

Note. This is a literate programming implementation directly linked to the application source code. Open this document at the command line, within its directory, i.e.,

```
@tarheellinux/hermes:texmacs doc.tm &
```

Links within a light green background [Makefile](#) directly invoke an editor to modify the specified file. Syntax-highlighted images of the file are automatically updated within this document.

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1 Theory

Compute the force exerted by escaping gas in a 1D model of the *Hermes* spaceship. Solve 2D Euler equations in conservative form

$$q_t + f(q)_x + g(q)_y = 0$$

$$q = \begin{pmatrix} \rho \\ \rho u \\ \rho v \\ \rho E \end{pmatrix}, f(q) = \begin{pmatrix} \rho u \\ \rho u^2 + p \\ \rho u v \\ \rho u H \end{pmatrix}, g(q) = \begin{pmatrix} \rho v \\ \rho u v \\ \rho v^2 + p \\ \rho v H \end{pmatrix},$$

through wave propagation algorithm using a linearized Riemann solver based on Roe averages. Set initial conditions to have no variation along y axis. Assume *Hermes* atmosphere is a perfect gas satisfying thermodynamic relations

$$p = \rho RT = \rho(\gamma - 1)c_v T = (\gamma - 1)\rho e = (\gamma - 1)\rho \left(E - \frac{u^2 + v^2}{2} \right),$$

$$H = E + \frac{p}{\rho}.$$

The sound speed is defined as $c^2 = \gamma p / \rho$.

1.1 Euler equation eigenmodes

The eigenmodes of the flux Jacobian matrices

$$A = \frac{\partial f}{\partial q}, B = \frac{\partial g}{\partial q},$$

are required for the Riemann solver. Define the conservative variables and fluxes.

Mathematica

```
In[33]:= q={rho, l, m, epsilon}
```

```
{rho, l, m, epsilon}
```

```
In[34]:= p=(gamma-1)(epsilon - (l^2+m^2)/(2 rho))
```

$$(\gamma - 1) \left(\epsilon - \frac{l^2 + m^2}{2\rho} \right)$$

```
In[35]:= H=(epsilon+p)/rho
```

$$\frac{(\gamma - 1) \left(\epsilon - \frac{l^2 + m^2}{2\rho} \right) + \epsilon}{\rho}$$

```
In[36]:= f={l, l^2/rho+p, l m/rho, l H}
```

$$\left\{ l, (\gamma - 1) \left(\epsilon - \frac{l^2 + m^2}{2\rho} \right) + \frac{l^2}{\rho}, \frac{l m}{\rho}, \frac{l \left((\gamma - 1) \left(\epsilon - \frac{l^2 + m^2}{2\rho} \right) + \epsilon \right)}{\rho} \right\}$$

```
In[37]:= g={m, l m/rho, m^2/rho+p, m H}
```

$$\left\{ m, \frac{l m}{\rho}, (\gamma - 1) \left(\epsilon - \frac{l^2 + m^2}{2\rho} \right) + \frac{m^2}{\rho}, \frac{m \left((\gamma - 1) \left(\epsilon - \frac{l^2 + m^2}{2\rho} \right) + \epsilon \right)}{\rho} \right\}$$

```
In[38]:=
```

Compute the Jacobians

In[38] := A=Simplify[Grad[f,q]]

$$\begin{pmatrix} 0 & 1 & 0 & 0 \\ \frac{(\gamma-3)l^2+(\gamma-1)m^2}{2\rho^2} & -\frac{(\gamma-3)l}{\rho} & \frac{m-\gamma m}{\rho} & \gamma-1 \\ -\frac{lm}{\rho^2} & \frac{m}{\rho} & \frac{l}{\rho} & 0 \\ \frac{l((\gamma-1)l^2+(\gamma-1)m^2-\epsilon\gamma\rho)}{\rho^3} & -\frac{3(\gamma-1)l^2+(\gamma-1)m^2-2\epsilon\gamma\rho}{2\rho^2} & -\frac{(\gamma-1)lm}{\rho^2} & \frac{\gamma l}{\rho} \end{pmatrix}$$

In[39] := B=Simplify[Grad[g,q]]

$$\begin{pmatrix} 0 & 0 & 1 & 0 \\ -\frac{lm}{\rho^2} & \frac{m}{\rho} & \frac{l}{\rho} & 0 \\ \frac{(\gamma-1)l^2+(\gamma-3)m^2}{2\rho^2} & \frac{l-\langle\text{gammal}\rangle}{\rho} & -\frac{(\gamma-3)m}{\rho} & \gamma-1 \\ \frac{m((\gamma-1)l^2+(\gamma-1)m^2-\epsilon\gamma\rho)}{\rho^3} & -\frac{(\gamma-1)lm}{\rho^2} & -\frac{(\gamma-1)l^2+3(\gamma-1)m^2-2\epsilon\gamma\rho}{2\rho^2} & \frac{\gamma m}{\rho} \end{pmatrix}$$

In[40] :=

Physical interpretation of the eigenmodes is easier in the primitive variables $\{\rho, u, v, p\}$. Define the transformation rules, and apply them to the flux Jacobian

In[23] := q2p = {l->rho u, m->rho v,
epsilon->rho c^2/gamma/(gamma-1)+rho(u^2/2+v^2/2)}

$$\left\{ l \rightarrow \rho u, m \rightarrow \rho v, \epsilon \rightarrow \frac{c^2 \rho}{(\gamma-1)\gamma} + \rho \left(\frac{u^2}{2} + \frac{v^2}{2} \right) \right\}$$

In[41] := Au=FullSimplify[A /. q2p]

$$\begin{pmatrix} 0 & 1 & 0 & 0 \\ \frac{1}{2}((\gamma-3)u^2+(\gamma-1)v^2) & -(\gamma-3)u & v-\gamma v & \gamma-1 \\ -uv & v & u & 0 \\ \frac{u((\gamma-2)(\gamma-1)(u^2+v^2)-2c^2)}{2(\gamma-1)} & \frac{2c^2-(\gamma-1)((2\gamma-3)u^2-v^2)}{2(\gamma-1)} & -(\gamma-1)uv & \gamma u \end{pmatrix}$$

In[42] := Bu=FullSimplify[B /. q2p]

$$\begin{pmatrix} 0 & 0 & 1 & 0 \\ -uv & v & u & 0 \\ \frac{1}{2}((\gamma-1)u^2+(\gamma-3)v^2) & u-\langle\text{gammau}\rangle & -(\gamma-3)v & \gamma-1 \\ \frac{v((\gamma-2)(\gamma-1)(u^2+v^2)-2c^2)}{2(\gamma-1)} & -(\gamma-1)uv & \frac{2c^2+(\gamma-1)(u^2+(3-2\gamma)v^2)}{2(\gamma-1)} & \gamma v \end{pmatrix}$$

In[43] :=

Disturbances in the fluid field values propagate as waves. The eigenvalues of the flux Jacobians are wave speeds

In[45] := lambdaA=FullSimplify[Eigenvalues[Au] /. q2p, Assumptions->{rho>0, c>0}]

$\{c+u, u-c, u, u\}$

`In[46] := lambdaB=FullSimplify[Eigenvalues[Bu] /. q2p, Assumptions->{rho>0, c>0}]`

$\{c+v, v-c, v, v\}$

`In[47] :=`

The Euler system has 4 eigenmodes with speeds $\lambda_1 = u + c$, $\lambda_2 = u - c$, $\lambda_3 = u$, $\lambda_4 = u$. Modes 1 and 2 correspond to forward and backward propagating acoustic waves. Modes 3 and 4 correspond to the same eigenvalue and correspond to propagation of a shear wave and a contact discontinuity.

The eigenvectors of the flux Jacobians describe how the conservative variables are coupled in wave propagation. After computation of the eigenvectors, the modes are expressed in a simpler, physically meaningful form using allowable algebraic operations.

`In[63] := RA=FullSimplify[Eigenvectors[Au] /. q2p, Assumptions->{rho>0, c>0}];
Transpose[RA]`

$$\begin{pmatrix} \frac{2(\gamma-1)}{2c^2+2(\gamma-1)uc+(\gamma-1)(u^2+v^2)} & \frac{2(\gamma-1)}{2c^2-2(\gamma-1)uc+(\gamma-1)(u^2+v^2)} & \frac{2}{u^2-v^2} & \frac{2v}{v^2-u^2} \\ \frac{2(\gamma-1)(c+u)}{2c^2+2(\gamma-1)uc+(\gamma-1)(u^2+v^2)} & \frac{2(\gamma-1)(c-u)}{2c^2-2(\gamma-1)uc+(\gamma-1)(u^2+v^2)} & \frac{2u}{u^2-v^2} & -\frac{2uv}{u^2-v^2} \\ \frac{2(\gamma-1)v}{2c^2+2(\gamma-1)uc+(\gamma-1)(u^2+v^2)} & \frac{2(\gamma-1)v}{2c^2-2(\gamma-1)uc+(\gamma-1)(u^2+v^2)} & 0 & 1 \\ 1 & 1 & 1 & 0 \end{pmatrix}$$

`In[64] := RA[[1]]=Simplify[RA[[1]]/RA[[1,1]] /. c^2 -> (gamma-1)(h-(u^2+v^2)/2)];
RA[[2]]=Simplify[RA[[2]]/RA[[2,1]] /. c^2 -> (gamma-1)(h-(u^2+v^2)/2)];
RA[[4]]=Simplify[RA[[4]]/RA[[3,1]]];
RA[[3]]=Simplify[RA[[3]]/RA[[3,1]]];
Transpose[RA]`

$$\begin{pmatrix} 1 & 1 & 1 & -v \\ c+u & u-c & u & -uv \\ v & v & 0 & \frac{1}{2}(u^2-v^2) \\ h+cu & h-cu & \frac{1}{2}(u^2-v^2) & 0 \end{pmatrix}$$

`In[65] := RA[[4]]=Simplify[RA[[4]] + v RA[[3]]]; Transpose[RA]`

$$\begin{pmatrix} 1 & 1 & 1 & 0 \\ c+u & u-c & u & 0 \\ v & v & 0 & \frac{1}{2}(u^2-v^2) \\ h+cu & h-cu & \frac{1}{2}(u^2-v^2) & \frac{1}{2}v(u^2-v^2) \end{pmatrix}$$

`In[66] := RA[[4]]=Simplify[RA[[4]]/RA[[4,3]]]; Transpose[RA]`

$$\begin{pmatrix} 1 & 1 & 1 & 0 \\ c+u & u-c & u & 0 \\ v & v & 0 & 1 \\ h+cu & h-cu & \frac{1}{2}(u^2-v^2) & v \end{pmatrix}$$

`In[67] := RA[[3]]=FullSimplify[RA[[3]] + v RA[[4]]]; Transpose[RA]`

$$\begin{pmatrix} 1 & 1 & 1 & 0 \\ c+u & u-c & u & 0 \\ v & v & v & 1 \\ h+cu & h-cu & \frac{1}{2}(u^2+v^2) & v \end{pmatrix}$$

In[70] := XA=RA; Transpose[XA]

$$\begin{pmatrix} 1 & 1 & 1 & 0 \\ c+u & u-c & u & 0 \\ v & v & v & 1 \\ h+cu & h-cu & \frac{1}{2}(u^2+v^2) & v \end{pmatrix}$$

In[71] :=

2 Makefile

The `Makefile` contains application-specific instructions to produce an `xbear` executable.

```

=====
# BEARCLAW Boundary Embedded Adaptive Refinement Conservation LAW package
=====
# (c) Copyright Sorin Mitran, 2017
# Department of Mathematics
# University of North Carolina at Chapel Hill
# mitran@unc.edu
# -----
# This code may be freely used for educational and research purposes.
# For any other use please contact the author.
# -----
# File:      Makefile
# Purpose:   Build 2D Euler equation Hermes breach model
# Contains:
# Revision History: Ver. 1.0 Feb. 2017 Sorin Mitran
# -----
#
default: xbear

include $(BEARCLAW)/Makefile.inc

PROBLEM_SOURCES = \
  nodeinfodef.f90 \
  problem.f90

PROBLEM_OBJECTS = \
  nodeinfodef.o \
  problem.o

METHOD_SOURCES = \
  $(BEARCLAWLIB)/wavebear.f90

METHOD_MODULES = \
  $(BEARCLAWLIB)/wavebear.mod

METHOD_OBJECTS = \
  $(BEARCLAWLIB)/wavebear.o

nodeinfodef.o: nodeinfodef.f90
$(FC) -$(DIR_MODULES) -o nodeinfodef.o nodeinfodef.f90

problem.o: problem.f90 physflux1.f90
$(FC) -$(DIR_MODULES) -o problem.o problem.f90

include $(BEARCLAW)/Makefile.targets

```

3 Input files for Hermes 2D model with multiple root-level grids

Multiple root-level grids are defined in order to model the *Hermes* substructures.

The number of BEARCLAW root level grids is specified in `bear.data`. Each grid is defined through a `grid nnn .data` file. The first grid is defined in `grid.data` (the most common case), additional grids are defined by `grid nnn .data` with nnn starting from 002 to the number of grids. The relative positions of each grid have to be specified; this is done through a global numbering system (i, j) .

3.1 Problem dependent parameters

Define in `setprob.data` the state inside the *Hermes*, and that approximating vacuum conditions. The data in this file is read in by `problem.f90:setprob` and is typically stored as global variables defined in `problem.f90:GlobalVariables`.

```
T          Entropy fix (efix)
1.4d0      gamma
0.         xBreach
42390      mHermes (kg)
1.0130000000000000d5 1.2250000000000000d0 0.0000000000000000d0 0.0000000000000000d0 Hermes atmosphere
0.0100000000000000d5 0.0456000000000000d0 0.0000000000000000d0 0.0000000000000000d0 Vacuum approximation
p (Pa)      rho (kg/m^3)    u (m/s)      v (m/s)
```

Mass calculation:

```
Hermes is a L=100m long x D=5m diameter cylinder with d=5mm thick walls made of Al, rhoAl=2700 kg/m^3
mHull = rhoAl*L*(pi*D)*d = 21195 kg
mPayload = mHull

mHermes ~= 42390 kg
```

3.2 Global run parameters

The `bear.data` file contains two sections.

1. `bear.data:RunFlags` sets flags affecting various global execution options

```
!:RunFlags:! | Variable | Description
=====
F 0 Restart, Frame Resume from checkpoint data dump
F LevelEqSets Solve different equations on grid levels
F LevelMethods Apply different algorithms on grid levels
F SaveAtFixedTimes F=maintain CFL, T=save data at desired times
F MaintainAuxArrays Treat aux similarly to q in MPI runs
F InitialAMRonly Generate initial AMR structure and stop
T OutputStyleParams Outputstyle line contains additional formatting
=====
```

2. `bear.data:RunParameters` sets parameters affecting various global execution options

```
!:RunParameters:!
=====
14 nRootGrids Number of root-level grids
3 MaxLevels Maximum number of grid refinement levels
2 2 2 2 2 CoarsenRatio ... of child grid to obtain parent spacing
4 MinimumGridPoints ... along one dimension
1 TimeStepMethod 0 fixed dt, 1 variable dt
0.d0 t0 initial time (if not Restart)
1.5d0 tfinal final time
4 0.5 MaxCFLRetry, rCFL Try reducing CFL by this ratio this many times
3 OutputStyle 1 AMRCLAW, 2 TECPLOT, 3 HDF, 9 GnuPlot, 11 VTK
300 OutputFrames Number of data checkpoints
T T T T T T OutputLevel Level output flag
=====
```

3.3 Root-level grid parameters

Define the individual grids with minimal effort through use of the Python template facility.

1. `grid.template:GridParameters` defines the grid geometry and boundaries. Variables that change with each grid are prefaced by a dollar sign.

```
!:GridParameters:! Variable Description
=====
2 nDim Grid spatial dimensions
4 MaxLevel Max grid refinement levels for this grid
500 mx Cells in x direction
4 my Cells in y direction
1 100 mGlobal(1) Global index extents of this grid (x-direction)
40 50 mGlobal(2) Global index extents of this grid (y-direction)
-10.d0 xlower Left edge of computational domain
0.d0 xupper Right edge of computational domain
0.0d0 ylower Bottom edge of computational domain
4.0d0 yupper Top edge of computational domain
2 mbc Number of ghost cells at each boundary
1 mthbc(1) Left boundary condition code
11 mthbc(2) Right boundary condition code
3 mthbc(3) Bottom boundary condition code
3 mthbc(4) Top boundary condition code
0.15d-3 dtv(1) Initial time step (constant dt TimeStepMethod=0)
1.0d99 dtv(2) Max allowable time step
1.00d0 cflv(1) Max allowable Courant number
0.90d0 cflv(2) Desired Courant number
0.1 cflv(3) Time step relaxation parameter
=====
```

2. `grid.template:MultiphysicsParameters` defines the conservation laws solved on this grid. These are identical for all grids in the computation

```
!MultiphysicsParameters: ! one value if LevelEqSets==F else (>=MaxLevel) values
!=====
! NrVars      = Number of primary field variables
4
! Output style parameters
0 1 1 0
! nEquationSet = Equation set for these fields
1
! maux        = Number of auxilliary fields
0
!=====
```

3. `grid.template:NumericalSchemeParameters` defines the numerical scheme on this grid, same for all grids in the computation.

```
!NumericalSchemeParameters: ! one value if LevelMethods==F else (>=MaxLevel) values
!=====
0      method(1) = (reserved)
2      method(2) = convergence order
2      method(3) = transverse convergence order
0      method(4) = verbosity of wavebear output
0      method(5) = source term splitting
0      method(6) = 0 split q differences, 1 split flux differences
0      method(7) = radius of slab around current 1D array of cells

4      mwaves    = number of waves in each Riemann solution
3 3 3 3      mthlim(mw) = limiter for each wave (mw=1,mwaves)
!=====
```

4. `grid.template>UserRootLevelParameters` defines additional grid parameters

```
!UserRootLevelParameters: !
!=====
! (none for this application)
!=====
```

Define arrays for the variables that reflect geometry in Fig. 1. It's convenient to define a grid number array to keep visual track of grid dimension assignment.

```
Python] nG      =[ 1, 2, 3, 4, 5, 6, 7, 8, 9, 10, 11, 12, 13, 14];
Python] mx      =[100,200,200,250,150, 50,200, 50, 50, 50,150, 50,200,500];
Python] my      =[100, 50,100,100,200,200,200,150,100,150,200,200,200,100];
Python] mGx1    =[ 1,101,101,301,401,551,601,551,551,551,401,551,601,601];
Python] mGx2    =[100,300,300,550,550,600,800,600,600,600,550,600,800,1100];
Python] mGy1    =[351,451,351,351,501,501,501,451,351,201, 1, 1, 1,351];
Python] mGy2    =[450,500,450,450,800,800,800,600,450,350,200,200,200,450];
Python] xlower=[ 0, 10, 10, 30, 40, 55, 60, 55, 55, 55, 40, 55, 60, 60];
Python] xupper=[ 10, 30, 30, 55, 55, 60, 80, 60, 60, 60, 55, 60, 80,110];
Python] ylower=[ 35, 45, 35, 35, 60, 60, 60, 45, 35, 20, 0, 0, 0, 35];
Python] yupper=[ 45, 50, 45, 45, 80, 80, 80, 60, 45, 35, 20, 20, 20, 45];
```



```
Python] mthbc1=[ 1, 10,999,999, 10,999,999, 10,999, 10, 10,999,999,999];
Python] mthbc2=[999, 11,999,999,999,999, 11, 11,999, 11,999,999, 11, 11];
Python] mthbc3=[ 3,999, 3, 3, 3,999, 3,999,999,999, 3, 3, 3, 3];
Python] mthbc4=[ 3, 3,999, 3, 3, 3, 3,999,999,999, 3,999, 3, 3];
Python] nG    =[ 1, 2, 3, 4, 5, 6, 7, 8, 9, 10, 11, 12, 13, 14];
          nGrids=len(nG); print nGrids
```

14

```
Python]
```

Use the Python template module.

```
Python] from string import Template;
```

Write a loop to replace template variables with values defined above and write resulting grid*nnn*.data files.

```
Python] s=open("grid.template","r").read(); t=Template(s);
for i in range(nGrids):
    g=t.safe_substitute(mx=str(mx[i]),my=str(my[i]),mGx1=str(mGx1[i]),
mGx2=str(mGx2[i]),mGy1=str(mGy1[i]),mGy2=str(mGy2[i]),xlower=str(xlower[i]),
xupper=str(xupper[i]),ylower=str(ylower[i]),yupper=str(yupper[i]),
mthbc1=str(mthbc1[i]),mthbc2=str(mthbc2[i]),mthbc3=str(mthbc3[i]),
mthbc4=str(mthbc4[i]) );
    i1=i+1; fname = 'grid%3d.data' % i1
    fgrid = open(fname, 'w');
    fgrid.write(g);
    fgrid.close();
```

```
Python]
```

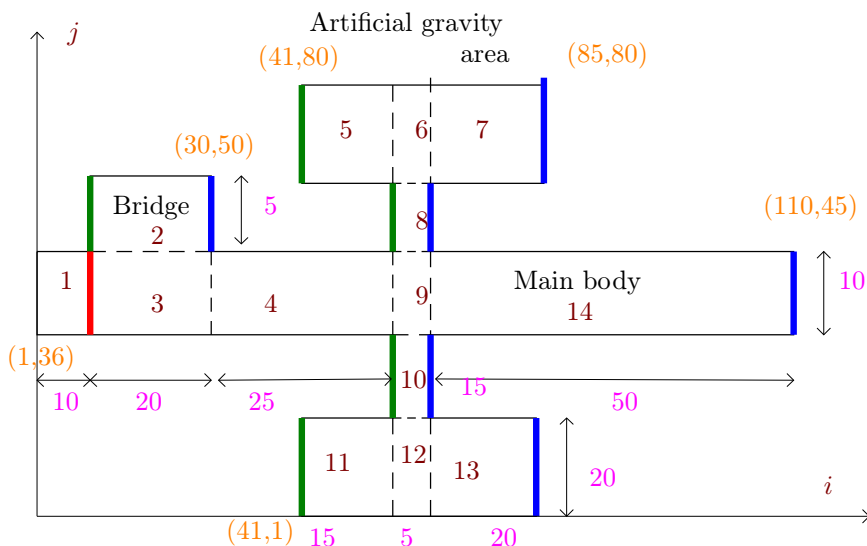


Figure 1. Hermes schematic model showing root level grids and boundary conditions: (red) breach location, (green) forward moving wall, (blue) rear moving wall. Dimensions are shown in meters (magenta). Global indexing coordinate pairs of cell centers (assuming cell size of 1m) are shown in orange.

4 Problem definition module

The `problem.f90` file defines a Fortran module that encapsulates problem physics.

4.1 Global variables

`problem.f90:GlobalVariables` specifies variables globally defined throughout the module

```
!GlobalVariables:! Problem specific global variables
! Parameters
INTEGER, PARAMETER :: OK=0
REAL (KIND=qPrec), PARAMETER :: zero=0., half=0.5d0, one=1.d0, two=2.d0
! Variables
LOGICAL efix                ! Flag: efix=T to apply entropy fix
REAL (KIND=qPrec) :: gamma,gamma1    ! Adiabatic coefficient gamma, gamma-1
REAL (KIND=xPrec) :: xBreach          ! Breach location
REAL (KIND=qPrec) :: vHermes,mHermes  ! Hermes velocity, mass
REAL (KIND=qPrec), DIMENSION(2) :: pQ,rhoQ,uQ,vQ ! Initial values

!=====
CONTAINS
!=====
```

4.2 Set problem parameters

`problem.f90:setprob` reads problem-specific parameters

```
!setprob:! Read problem global data
SUBROUTINE setprob
  INTEGER i
  ! Read gas parameters from file
  OPEN(UNIT=7,FILE='setprob.data',STATUS='old',FORM='formatted')
  READ(7,*)efix                ! set to T to apply entropy fix in Riemann solver
  READ(7,*) gamma              ! adiabatic constant
  gamma1 = gamma - one
  READ(7,*) xBreach            ! position of Hermes breach
  READ(7,*) mHermes
  DO i=1,2
    READ(7,*)pQ(i),rhoQ(i),uQ(i),vQ(i) ! initial values (primitive variables)
  END DO
  CLOSE(7)
  ! Initialize Hermes velocity, reaction force of escaping gas
  vHermes=zero
END SUBROUTINE setprob
```

4.3 Field variable initialization

`problem.f90:qinit` initialized the field variables

```
!qinit:! Field variable initialization
SUBROUTINE qinit(Info)
  TYPE (NodeInfo) :: Info ! Data associated with this grid
  ! Internal declarations
  INTEGER i,iQ
  REAL (KIND=xPrec) :: x
  ! iQ=1 inside Hermes, at right; iQ=2 vacuum approximation
  x=Info%Xlower(1)+Info%dX(1)/2.0
  DO i=1,Info%MX(1)
    IF (x >= xBreach) THEN
      iQ = 1
    ELSE
      iQ = 2
    ENDIF
    Info%q(i,1,1,1) = rhoQ(iQ)
    Info%q(i,1,1,2) = rhoQ(iQ)*uQ(iQ)
    Info%q(i,1,1,3) = rhoQ(iQ)*vQ(iQ)
    Info%q(i,1,1,4) = pQ(iQ)/gamma1 + 0.5d0*rhoQ(iQ)*(uQ(iQ)**2 + vQ(iQ)**2)
    x=x+Info%dX(1)
  END DO
END SUBROUTINE qinit
```

4.4 Boundary conditions

`problem.f90:problemBC` defines application-specific boundary conditions, different from the predefined codes: 1=outflow, 2=periodic, 3=wall.

```
!problemBC:! User-defined boundary conditions
SUBROUTINE problemBC(Info)
  TYPE (NodeInfo) :: Info ! Data associated with this grid
  ! Internal declarations
  INTEGER ibc,mx,mbc,r,rmbc
  SELECT CASE(Info%bcnow)
  CASE (1)
    SELECT CASE(Info%thbc(1))
    CASE (10) ! Moving solid wall boundary condition at left
      mx=Info%MX(1); mbc=Info%mbc; r=Info%r; rmbc=r*mbc
      DO ibc=1,rmbc
        Info%q(1-ibc,1,1,1) = Info%q(ibc,1,1,1)
      END DO
      ! Negate the normal velocity (assumed to 2nd component):
      DO ibc=1,rmbc
        Info%q(1-ibc,1,1,2) = 2*vHermes - Info%q(ibc,1,1,2)
      END DO
    END SELECT
  CASE (2)
    SELECT CASE(Info%thbc(2))
    CASE (11) ! Moving solid wall boundary condition at right
      mx=Info%MX(1); mbc=Info%mbc; r=Info%r; rmbc=r*mbc
      DO ibc=1,rmbc
        Info%q(mx+ibc,1,1,1) = Info%q(mx+1-ibc,1,1,1)
      END DO
      ! Negate the normal velocity (assumed to 2nd component):
      DO ibc=1,rmbc
        Info%q(mx+ibc,1,1,2) = 2*vHermes - Info%q(mx+1-ibc,1,1,2)
      END DO
    END SELECT
  END SELECT
END SUBROUTINE problemBC
```

4.5 Actions to take after each time step

The force acting on wall perpendicular to the x -direction is computed in `problem.f90:afterstep` and used to determine the change in *Hermes* velocity.

```

!afterstep: Actions after a time step
SUBROUTINE afterstep(Info)
  TYPE (NodeInfo) :: Info
  ! Internal declarations
  INTEGER j,mx,my
  REAL (KIND=qPrec) :: rho,rhou,rhov,rhoE,p,F
  ! Check if this is a leaf node grid with a moving wall boundary condition on right
  IF ((Info%SubGrids .EQ. 0) .AND. (Info%nthbc(2) .EQ. 10)) THEN
    ! Yes, moving wall at left edge. Compute force, update velocity
    my = Info%mx(2); F=0.
    DO j=1,my
      rho = Info%q(1,j,1,1);
      rhou = Info%q(1,j,1,2);
      rhov = Info%q(1,j,1,3);
      rhoE = Info%q(1,j,1,4);
      p = gamma1*(rhoE - (rhou**2+rhov**2)/(2.*rho))
      F = F - p*Info%dX(2)
    ENDDO
    ! Increment change in Hermes velocity produced by each leaf grid
    vHermes = vHermes + Info%dt*F/mHermes
  ENDIF
  ! Check if this is a leaf node grid with a moving wall boundary condition on right
  IF ((Info%SubGrids .EQ. 0) .AND. (Info%nthbc(2) .EQ. 11)) THEN
    ! Yes, moving wall at right edge. Compute force, update velocity
    mx = Info%mx(1); my = Info%mx(2); F=0.
    DO j=1,my
      rho = Info%q(mx,j,1,1);
      rhou = Info%q(mx,j,1,2);
      rhov = Info%q(mx,j,1,3);
      rhoE = Info%q(mx,j,1,4);
      p = gamma1*(rhoE - (rhou**2+rhov**2)/(2.*rho))
      F = F + p*Info%dX(2)
    ENDDO
    ! Increment change in Hermes velocity produced by each leaf grid
    vHermes = vHermes + Info%dt*F/mHermes
  ENDIF
END SUBROUTINE afterstep

```

4.6 Physical fluxes

The `physflux1.f90:physflux` routine defines the problem physics through specification of the fluxes and solution of the Riemann problem

4.6.1 Local variable declarations

The quantities $\Delta q, \alpha$ that arise in the Riemann problem solution are stored in `delta`, `coef` defined in `physflux1.f90:physfluxInternalDeclarations`. These arrays are allocated as needed. The `SAVE` attribute maintains the values in `delta`, `coef` between successive request calls to `physflux`.

```
!physfluxInternalDeclarations:
INTEGER, SAVE :: mbc,mx,nq,mwav,NrVars,mwaves,iError,j,mu,mv,mw,i
INTEGER, POINTER, DIMENSION(:) :: iCell,iEdge,iLft,iRgt
REAL (KIND=qPrec), ALLOCATABLE, SAVE, DIMENSION(:) :: delta,coef
! Cell center quantities
REAL (KIND=qPrec), ALLOCATABLE, SAVE, DIMENSION(:) :: sqrtrho,pres, &
    u,v,c,h,rho
! Roe averaged quantities
REAL (KIND=qPrec), ALLOCATABLE, SAVE, DIMENSION(:) :: uR,vR,hR,cR, &
    rhosq2,g1c2,u2v2,euv
! Shorter local names
REAL (KIND=qPrec), POINTER, DIMENSION(,:) :: Apdq,Amdq,Asdq,BpAsdq,BmAsdq, &
    speed,q1D
REAL (KIND=qPrec), POINTER, DIMENSION(,,:) :: wave
REAL (KIND=qPrec) :: rho1,rhou1,en1,p1,c1,rho2,rhou2,en2,p2,c2
REAL (KIND=qPrec) :: s0,s1,s2,s3,sfract,df(5),rhov1,rhov2
```

4.6.2 Local aliases for Info fields

It is convenient to have shorter names for the `Info` structure fields, as in `wave=>Info%wave` - the local `wave` pointer references the `Info%wave` array. Define these pointers in `physflux1.f90:physfluxInfo`

```
!physfluxInfo: Extract information from Info structure
! Current number of interior cells, boundary cells, field variables, waves
mx=Info%mxnow; mbc=Info%mbc; NrVars=Info%NrVars; mwaves=Info%mwaves
! Associate local names with Info components
q1D=>Info%q1D
Apdq=>Info%Apdq; Amdq=>Info%Amdq; speed=>Info%speed; wave=>Info%wave
Asdq=>Info%Asdq; BpAsdq=>Info%BpAsdq; BmAsdq=>Info%BmAsdq
iEdge=>Info%i1D; iCell=>Info%i1Dcells; iLft=>Info%i1Dleft; iRgt=>Info%i1Dright
```

4.6.3 Switching between Riemann problem directions

The eigenmodes for $A = \partial f / \partial q$ and $B = \partial g / \partial q$ can be obtained from one another through index permutation as defined in `physflux1.f90:physfluxNormalDirection`.

```
!physfluxNormalDirection: Define permutation to reuse eigenmode code
IF (ixy==1) THEN
    mu = 2; mv = 3
ELSE
    mu = 3; mv = 2
END IF
SELECT CASE (irequest)
```

4.6.4 Riemann problem solution

Code must be written to solve the normal Riemann problem along one-dimensional slices of the grid in response to `RequestNormalWaves`, and the transverse problem in response to `RequestTransverseWaves`.

The `physflux` routine is also called with requests `Initialize`, `Finalize`, `RequestFluxes` by the `wavebear` solver module.

Initialize The code in `physflux1.f90:physfluxInitialize` allocates space needed for cell-centered quantities (`u`, `v`, `h`, `c`), Roe-average quantities (`uR`, `vR`, `hR`, `cR`), and various work vectors.

```
!physfluxInitialize!
CASE (Initialize)
! Allocate local work space
ALLOCATE(delta(NrVars),coef(mwaves),
          sqrtrho(1-mbc:mx+mbc),pres(1-mbc:mx+mbc),rho(1-mbc:mx+mbc),
          u(1-mbc:mx+mbc),v(1-mbc:mx+mbc),h(1-mbc:mx+mbc),c(1-mbc:mx+mbc),
          uR(2-mbc:mx+mbc),vR(2-mbc:mx+mbc),hR(2-mbc:mx+mbc),cR(2-mbc:mx+mbc),
          rhosq2(2-mbc:mx+mbc),g1c2(2-mbc:mx+mbc),u2v2(2-mbc:mx+mbc),
          euv(2-mbc:mx+mbc),
          STAT=iError)
IF (iError/=OK) THEN
PRINT *,'Cannot allocate work space in problem module, physflux'
STOP
END IF
```

Finalize The code in `physflux1.f90:physfluxFinalize` releases space previously allocated by the response to `Initialize`.

```
!physfluxFinalize!
CASE (Finalize)
! DeAllocate local work space
DEALLOCATE(delta,coef,sqrtrho,pres,rho,u,v,h,c,uR,vR,hR,cR,rhosq2,g1c2,
            u2v2,euv,STAT=iError)
IF (iError/=OK) THEN
PRINT *,'Cannot deallocate work space in problem module, physflux'
STOP
END IF
```

RequestFluxes Code in `physflux1.f90:physfluxRequestFluxes` evaluate the physical fluxes based on cell-centered quantities.

```
!physfluxRequestFluxes!
CASE (RequestFluxes)
IF (MINVAL(q1D(1-mbc:mx+mbc,1))<=zero) THEN
PRINT *,'Error: Negative or zero density in physflux.'
STOP
END IF
! Compute cell centered quantities required in Riemann solve & entropy fix
rho(iCell) = 1./q1D(iCell,1)
u(iCell) = q1D(iCell,mu)*rho(iCell)
v(iCell) = q1D(iCell,mv)*rho(iCell)
rho(iCell) = q1D(iCell,1)
sqrtrho(iCell) = SQRT(rho(iCell))
pres(iCell) = gamma1*( q1D(iCell,4) - &
                    half*(q1D(iCell,mu)*u(iCell) + q1D(iCell,mv)*v(iCell)) )
c = gamma*pres/rho
IF (MINVAL(c)<=zero) THEN
PRINT *,'Error: Non-physical centered sound velocity in physflux, RequestFluxes.'
STOP
END IF
c = SQRT(c)
h(iCell) = (q1D(iCell,4) + pres(iCell)) / rho(iCell)
! Compute the physical fluxes at cell centers
f(iCell,1) = q1D(iCell,mu)
f(iCell,mu) = f(iCell,1)*u(iCell)+pres(iCell)
f(iCell,mv) = f(iCell,1)*v(iCell)
f(iCell,4) = f(iCell,1)*h(iCell)
```

RequestNormalWaves The solution of the Riemann problem in the normal direction is computed in a succession of stages:

physflux1.f90:physfluxRequestNormalWaves Initialization of fluctuations and waves

```
!physfluxRequestNormalWaves:!  
CASE (RequestNormalWaves)  
  Apdq=0.; Amdq=0.; speed=0.; wave=0.
```

physflux1.f90:ComputeRoeAverages computes the intermediate state at which the flux Jacobians are evaluated.

```
!:ComputeRoeAverages!  
rhosq2(iEdge) = one / (sqrtrho(iLft)+sqrtrho(iRgt))  
uR(iEdge) = (q1D(iLft,mu)/sqrtrho(iLft) + q1D(iRgt,mu)/sqrtrho(iRgt)) * rhosq2(iEdge)  
vR(iEdge) = (q1D(iLft,mv)/sqrtrho(iLft) + q1D(iRgt,mv)/sqrtrho(iRgt)) * rhosq2(iEdge)  
hR(iEdge) = ((q1D(iLft,4)+pres(iLft))/sqrtrho(iLft) +  
             (q1D(iRgt,4)+pres(iRgt))/sqrtrho(iRgt)) * rhosq2(iEdge)  
u2v2=uR**2+vR**2  
cR = gamma1*(hR-half*u2v2)  
IF (MINVAL(cR)<=zero) THEN  
  PRINT *,'Error: Negative or zero Roe average sound velocity in physflux.'  
  STOP  
END IF  
g1c2 = gamma1/cR  
cR = SQRT(cR)  
euv = hR - u2v2
```

physflux1.f90:ConstructEigenmodes defines the propagating waves resulting from the flux Jacobian eigensystem

```
!:ConstructEigenmodes!  
! Backward acoustic  
speed(iEdge,1) = uR(iEdge) - cR(iEdge)  
wave(iEdge,1,1) = one  
wave(iEdge,mu,1) = speed(iEdge,1)  
wave(iEdge,mv,1) = vR(iEdge)  
wave(iEdge,4,1) = hR(iEdge) - uR(iEdge)*cR(iEdge)  
! Shear  
speed(iEdge,2) = uR(iEdge)  
wave(iEdge,1,2) = zero  
wave(iEdge,mu,2) = zero  
wave(iEdge,mv,2) = one  
wave(iEdge,4,2) = vR(iEdge)  
! Entropy  
speed(iEdge,3) = uR(iEdge)  
wave(iEdge,1,3) = one  
wave(iEdge,mu,3) = uR(iEdge)  
wave(iEdge,mv,3) = vR(iEdge)  
wave(iEdge,4,3) = half*u2v2(iEdge)  
! Forward acoustic  
speed(iEdge,4) = uR(iEdge) + cR(iEdge)  
wave(iEdge,1,4) = one  
wave(iEdge,mu,4) = speed(iEdge,4)  
wave(iEdge,mv,4) = vR(iEdge)  
wave(iEdge,4,4) = hR(iEdge) + uR(iEdge)*cR(iEdge)  
! Tracer  
IF (mwaves==5) THEN  
  speed(iEdge,5) = uR(iEdge)  
  wave(iEdge,1:4,5) = zero  
  wave(iEdge,5,5) = one  
END IF
```

`physflux1.f90:DecomposeJump` solves the linearized Riemann problem at each cell interface

```

!DecomposeJump:! in q onto eigenbases
DO j=2-mbc,mx+mbc
  ! Find coef(1) thru coef(4), the coefficients of the 4 eigenvectors
  delta(1) = q1D(j,1) - q1D(j-1,1)
  delta(2) = q1D(j,mu) - q1D(j-1,mu)
  delta(3) = q1D(j,mv) - q1D(j-1,mv)
  delta(4) = q1D(j,4) - q1D(j-1,4)
  coef(3) = g1c2(j) * (euv(j)*delta(1) + uR(j)*delta(2) + vR(j)*delta(3) - delta(4))
  coef(2) = delta(3) - vR(j)*delta(1)
  coef(4) = (delta(2) + (cR(j)-uR(j))*delta(1) - cR(j)*coef(3)) / (two*cR(j))
  coef(1) = delta(1) - coef(3) - coef(4)
  IF (NrVars==5) THEN
    ! Tracer variable
    coef(5) = q1D(j,5) - q1D(j-1,5)
  END IF
  DO mw=1,mwaves
    wave(j,1:NrVars,mw) = coef(mw)*wave(j,1:NrVars,mw)
  END DO
END DO

```

`physflux1.f90:GodunovFlux` evaluates the numerical flux at the interface between two finite volume cells $i, i + 1$ as the physical flux evaluated at the intermediate state

$$F_i^n = f(q_{i+1/2}^\uparrow)$$

with $q_{i+1/2}^\uparrow$ the solution to the Riemann problem at the $x_{i+1/2}$ interface. Through linearization of the Riemann problem, a continuous rarefaction fan has been replaced by a propagating rarefaction discontinuity that violates the thermodynamic condition of non-decreasing entropy. Only part of the jump associated with the rarefaction shock actually affects $q_{i+1/2}^\uparrow$. A sequence of checks is required to compute this correction (the “entropy fix”). Since this is computationally costly and only encountered in a small part of the finite volume cells within the computation, the entropy fix is an option specified in `setprob.data` through the logical flag `efix`.

```

!GodunovFlux:! for first-order approximation
IF (efix) THEN

```


`physflux1.f90:EntropyFix` checks each wave to determine if only a fraction of the decomposed jump should actually affect cells to the left.

```

!EntropyFix:
! Compute flux differences amdq and apdq.
! First compute amdq as sum of s*wave for left going waves.
! Incorporate entropy fix by adding a modified fraction of wave
! if s should change sign.
DO j=2-mbc,mx+mbc
! 1-wave. Check for fully supersonic case
s0=u(j-1)-c(j-1)
IF ( s0>=zero .AND. speed(j,1) > zero ) THEN
! Everything is right-going
Amdq(j,1:NrVars) = zero
CYCLE
END IF
! u-c to right of 1-wave
rho1 = q1D(j-1,1) + wave(j,1,1)
rho1 = q1D(j-1,mu) + wave(j,mu,1)
rho1 = q1D(j-1,mv) + wave(j,mv,1)
en1 = q1D(j-1,4) + wave(j,4,1)
p1 = gamma1*(en1 - 0.5*(rho1**2+rho1**2)/rho1)
IF ((p1<zero) .OR. (rho1<zero)) THEN
PRINT *, 'Negative pressure/density between 1-wave and 2,3-waves'
CYCLE
STOP
END IF
c1 = SQRT(gamma*p1/rho1)
s1 = rho1/rho1 - c1
IF ( s0<zero .AND. s1>zero ) THEN
! Transonic rarefaction in the 1-wave
sfract = s0 * (s1-speed(j,1)) / (s1-s0)
ELSE IF (speed(j,1) < zero) THEN
! Left-going 1-wave
sfract=speed(j,1)
ELSE
! Right-going 1-wave
sfract=zero ! Never should reach this instruction
END IF
Amdq(j,1:NrVars) = sfract*wave(j,1:NrVars,1)

! Check 2,3-wave (contact+shear discontinuities)
IF (speed(j,2) >= zero) CYCLE ! 2- and 3- wave are right-going
Amdq(j,1:NrVars) = Amdq(j,1:NrVars) + &
uR(j)*(wave(j,1:NrVars,2) + wave(j,1:NrVars,3))
! Check 4-wave. Compute u+c to left of 4-wave
rho2 = q1D(j,1) - wave(j,1,4)
rho2 = q1D(j,mu) - wave(j,mu,4)
rho2 = q1D(j,mv) - wave(j,mv,4)
en2 = q1D(j,4) - wave(j,4,4)
p2 = gamma1*(en2 - 0.5d0*(rho2**2+rho2**2)/rho2)
IF ((p2<=zero) .OR. (rho2<=zero)) THEN
PRINT *, 'Negative pressure/density between 2,3-waves and 4-wave'
CYCLE
STOP
END IF
c2 = SQRT(gamma*p2/rho2)
s2 = rho2/rho2 + c2; s3=u(j)+c(j)
IF ( s2 < zero .AND. s3 > zero ) THEN
! Transonic rarefaction in the 4-wave
sfract = s2 * (s3-speed(j,4)) / (s3-s2)
ELSE IF (speed(j,4) < zero) THEN
! Left-going 4-wave
sfract = speed(j,4)
ELSE
! Right-going 4-wave
CYCLE
END IF
Amdq(j,1:NrVars) = Amdq(j,1:NrVars) + sfract*wave(j,1:NrVars,4)
END DO

```

`physflux1.f90:RightGoingFluxes` subsequently evaluates the correction in the right-going fluxes

```

!RightGoingFluxes:!
DO j=2-mbc,mx+mbc
  df=zero
  DO mwav=1,mwaves
    df(1:NrVars) = df(1:NrVars) + speed(j,mwav) * wave(j,1:NrVars,mwav)
  END DO
  Apdq(j,1:NrVars) = df(1:NrVars) - Amdq(j,1:NrVars)
END DO
!===
ELSE
!===

```

`physflux1.f90:NoEntropyFix` directly uses the jump decomposition provided by the linearized Riemann solver to evaluate fluctuations that update cell averages, without correcting for possible rarefaction fans (and thus allowing small violations of the entropy condition in the computation).

```

!NoEntropyFix:!
DO mw=1,mwaves
  DO j=2-mbc,mx+mbc
    IF (speed(j,mw) < zero) THEN
      Amdq(j,1:NrVars) = Amdq(j,1:NrVars) + &
        speed(j,mw) * wave(j,1:NrVars,mw)
    ELSE
      Apdq(j,1:NrVars) = Apdq(j,1:NrVars) + &
        speed(j,mw) * wave(j,1:NrVars,mw)
    END IF
  END DO
END DO
END IF

```

RequestTransverseWaves Part of the fluctuation computed in the normal direction Riemann problem actually travels to adjacent cells in the transverse direction. This component is computed in `physflux1.f90:physfluxRequestTransverseWaves`

```

!:physfluxRequestTransverseWaves:
CASE (RequestTransverseWaves)
  BpAsdq=0.; BmAsdq=0.; speed=0.; wave=0.
  ! Construct the transverse eigenbasis
  ! Backward acoustic
  speed(iEdge,1) = vR(iEdge) - cR(iEdge)
  wave(iEdge,1,1) = one
  wave(iEdge,mu,1) = uR(iEdge)
  wave(iEdge,mv,1) = speed(iEdge,1)
  wave(iEdge,4,1) = hR(iEdge) - vR(iEdge)*cR(iEdge)
  ! Shear
  speed(iEdge,2) = vR(iEdge)
  wave(iEdge,1,2) = zero
  wave(iEdge,mu,2) = one
  wave(iEdge,mv,2) = zero
  wave(iEdge,4,2) = uR(iEdge)
  ! Entropy
  speed(iEdge,3) = vR(iEdge)
  wave(iEdge,1,3) = one
  wave(iEdge,mu,3) = uR(iEdge)
  wave(iEdge,mv,3) = vR(iEdge)
  wave(iEdge,4,3) = half*u2v2(iEdge)
  ! Forward acoustic
  speed(iEdge,4) = vR(iEdge) + cR(iEdge)
  wave(iEdge,1,4) = one
  wave(iEdge,mu,4) = uR(iEdge)
  wave(iEdge,mv,4) = speed(iEdge,4)
  wave(iEdge,4,4) = hR(iEdge) + vR(iEdge)*cR(iEdge)
IF (mwaves==5) THEN
  ! Tracer
  speed(iEdge,5) = vR(iEdge)
  wave(iEdge,1:4,5) = zero
  wave(iEdge,5,5) = one
END IF
  ! Decompose fluctuation in Asdq onto eigenbasis
DO j=2-mbc,mx+mbc
  ! Find coef(1) thru coef(4), the coefficients of the 4 eigenvectors
  delta(1:NrVars) = Asdq(j,1:NrVars)
  coef(3) = g1c2(j) * (euv(j)*delta(1) + uR(j)*delta(mu) + vR(j)*delta(mv) - delta(4))
  coef(2) = delta(mu) - uR(j)*delta(1)
  coef(4) = (delta(mv) + (cR(j)-vR(j))*delta(1) - cR(j)*coef(3)) / (two*cR(j))
  coef(1) = delta(1) - coef(3) - coef(4)
  IF (NrVars==5) THEN
  ! Tracer variable
  coef(5) = delta(5)
  END IF
  DO mw=1,mwaves
    wave(j,1:NrVars,mw) = coef(mw)*wave(j,1:NrVars,mw)
  END DO
END DO
  ! Compute the transverse fluctuations
  DO mwav=1,mwaves
  DO j=2-mbc,mx+mbc
    IF (speed(j,mwav) < zero) THEN
      BmAsdq(j,1:NrVars) = BmAsdq(j,1:NrVars) + &
        speed(j,mwav) * wave(j,1:NrVars,mwav)
    ELSE
      BpAsdq(j,1:NrVars) = BpAsdq(j,1:NrVars) + &
        speed(j,mwav) * wave(j,1:NrVars,mwav)
    END IF
  END DO
END DO

```

5 Results

The code is compiled and executed in a separate terminal window.

```
Shell session inside TeXmacs pid = 2590
```

```
Shell] pwd
```

```
/home/student/courses/MATH762/homework/hw1solution/1dshocktube
```

```
Shell] make xbear
```

```
/usr/bin/gfortran -c -Dgfortran -fcray-pointer -O3 -g -DHDF -DNOMKL -I/home/  
student/bearclaw/lib -o problem.o problem.f90  
/usr/bin/gfortran nodeinfodef.o problem.o \  
/home/student/bearclaw/lib/treeops.o /home/student/bearclaw/lib/cutcell.o /home/  
student/bearclaw/lib/fixup.o /home/student/bearclaw/lib/beario.o /home/student/  
bearclaw/lib/wavebear.o /home/student/bearclaw/lib/linsolve.o /home/student/
```

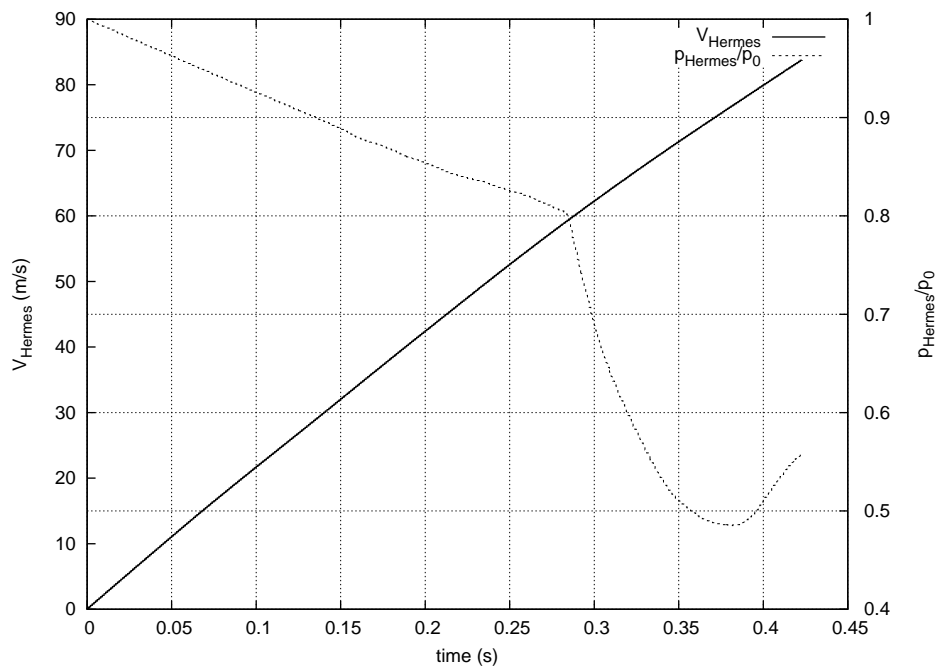
```
bearclaw/lib/infocfieldutils.o /home/student/bearclaw/lib/serial_exec.o /home/
student/bearclaw/lib/bearez.o /home/student/bearclaw/lib/setbc.o /home/student/
bearclaw/lib/driver.o -L/usr/lib64/hdf -lmfhdf -ldf -ljpeg -lz -L/usr/lib64 -
lhdf5_fortran -lhdf5 -o xbear
```

```
Shell] exo-open --launch TerminalEmulator xbear
```

```
Shell]
```

The code saves a time history of the *Hermes* velocity and pressure drop in `vhermes.data`, rendered below using Gnuplot

```
GNUplot] set grid x y2;
set xlabel("time (s)"); set ylabel("V_{Hermes} (m/s)");
set y2label("p_{Hermes}/p_0");
set autoscale y; set autoscale y2;
set ytics nomirror; set y2tics;
plot "vhermes.data" using 1:2 axes x1y1 title "V_{Hermes}" w l,
"vhermes.data" using 1:3 axes x1y2 title "p_{Hermes}/p_0" w l
```



```
GNUplot]
```

The OpenDX script `hermes.net` produces animations of the time evolution of the field variables. Choose File->Save Image ...->Continuous Saving to save the animation to `anim.miff`. Frames from the resulting animation can be extracted using `convert` and combined into a single image using `montage`.

```
Shell] rm --force anim.miff; exo-open --launch TerminalEmulator dx -execute -program
hermes.net
```

```
Shell] rm --force hermes???.png; convert anim.miff[0,10,30,45,60,80] hermes%03d.png
```

```
Shell] montage hermes???.png -tile 2x3 -geometry +0+0 hermes.png
```

```
Shell]
```

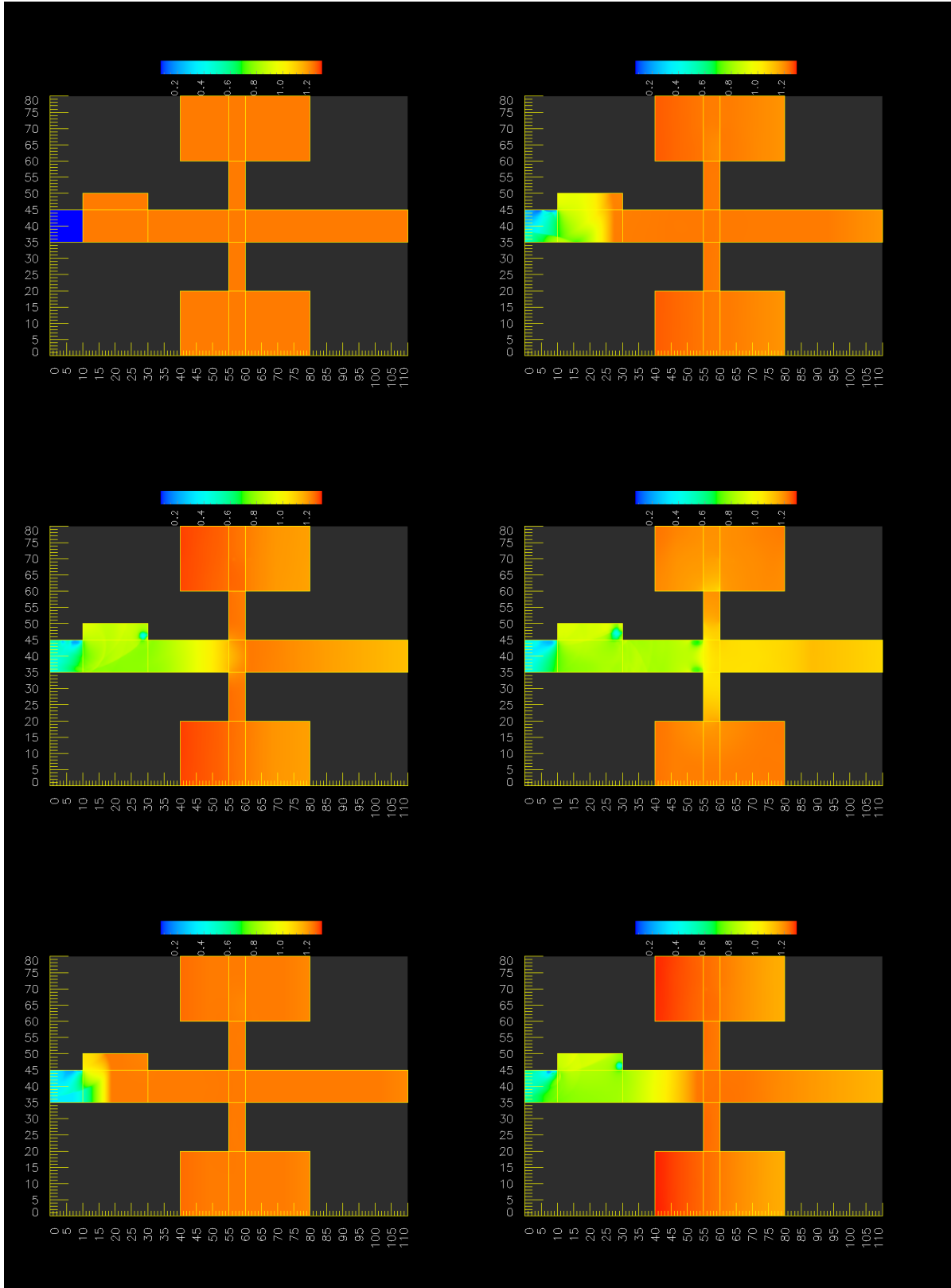


Figure 2. Evolution of density field with time in *Hermes*.

The OpenDX script `schlieren.net` produces an animation of $\log|\nabla\rho|$ as would be visualized using the Schlieren technique superimposed with the velocity variables. A Schlieren visualization shows the acoustic and vorticity waves in a compressible fluid.

```
Shell] rm --force schlieren.miff; exo-open --launch TerminalEmulator dx -execute -  
program schlieren.net
```

```
Shell] rm --force schlieren???.png; convert schlieren.miff[0,10,30,45,60,80]  
schlieren%03d.png
```

```
Shell] montage schlieren???.png -tile 2x3 -geometry +0+0 schlieren.png
```

```
Shell]
```

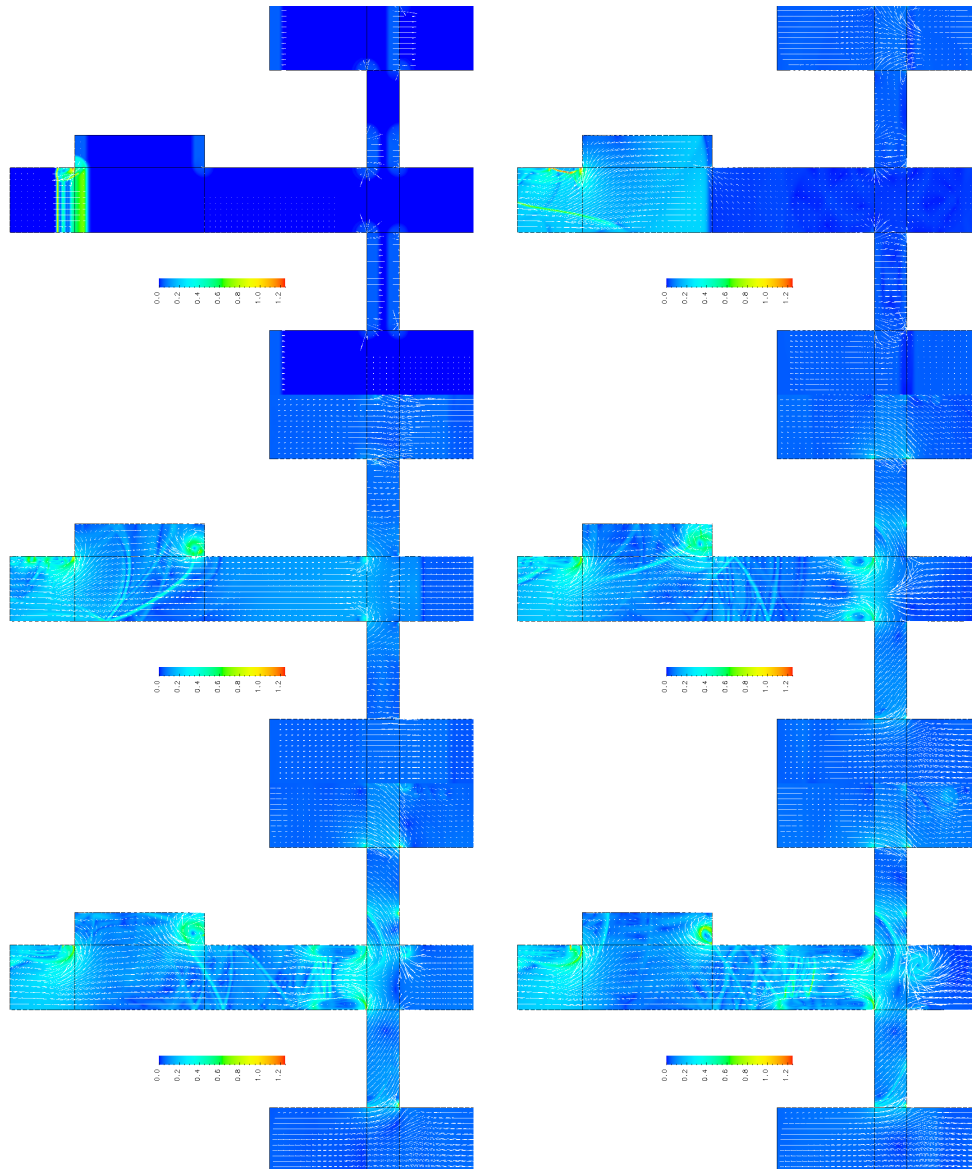


Figure 3. Schlieren visualization of *Hermes* breach.