



Overview

- Boltzmann equation, conservation laws
- Hermite polynomial expansions
- Equilibrium distribution



- $f(\mathbf{x}, \boldsymbol{\xi}, t)$ probability distribution function
- Force-free Boltzmann equation (dimensionless)

$$\frac{\partial f}{\partial t} + \xi_\alpha \frac{\partial f}{\partial x_\alpha} = \Omega(f)$$

- Conservation laws

Mass	$\int f(\mathbf{x}, \boldsymbol{\xi}, t) d\boldsymbol{\xi} =$	$\int f^{\text{eq}}(\mathbf{x}, \boldsymbol{\xi}, t) d\boldsymbol{\xi} =$	$\rho(\mathbf{x}, t)$
Momentum	$\int f(\mathbf{x}, \boldsymbol{\xi}, t) \boldsymbol{\xi} d\boldsymbol{\xi} =$	$\int f^{\text{eq}}(\mathbf{x}, \boldsymbol{\xi}, t) \boldsymbol{\xi} d\boldsymbol{\xi} =$	$\rho \mathbf{u}(\mathbf{x}, t)$
Total energy	$\frac{1}{2} \int f(\mathbf{x}, \boldsymbol{\xi}, t) \boldsymbol{\xi} ^2 d\boldsymbol{\xi} =$	$\frac{1}{2} \int f^{\text{eq}}(\mathbf{x}, \boldsymbol{\xi}, t) \boldsymbol{\xi} ^2 d\boldsymbol{\xi} =$	$\rho E(\mathbf{x}, t)$
Internal energy	$\frac{1}{2} \int f(\mathbf{x}, \boldsymbol{\xi}, t) \boldsymbol{\xi} - \mathbf{u} ^2 d\boldsymbol{\xi} =$	$\frac{1}{2} \int f^{\text{eq}}(\mathbf{x}, \boldsymbol{\xi}, t) \boldsymbol{\xi} - \mathbf{u} ^2 d\boldsymbol{\xi} =$	$\rho e(\mathbf{x}, t)$

$$E = e + \frac{1}{2} |\mathbf{u}|^2$$



- Definition from generating (weight) function

$$\omega(x) = \frac{1}{\sqrt{2\pi}} \exp\left(-\frac{x^2}{2}\right), H^{(n)}(x) = (-1)^n \frac{1}{\omega(x)} \frac{d^n}{dx^n} \omega(x)$$

```
In[1] := omega[x_] = Exp[-x^2/2]/Sqrt[2Pi]
```

$$\frac{e^{-\frac{x^2}{2}}}{\sqrt{2\pi}}$$

```
In[2] := H[n_, x_] := Expand[(-1)^n/omega[x] D[omega[x], {x, n}]];
Table[H[n, x], {n, 0, 5}]
```

```
{1, x, x^2 - 1, x^3 - 3x, x^4 - 6x^2 + 3, x^5 - 10x^3 + 15x}
```

```
In[3] := HHd1=Table[Integrate[omega[x] H[i,x] H[j,x], {x,-Infinity,Infinity}],
{ i, 0, 4}, {j, 0, 4}] // MatrixForm
```

$$\begin{pmatrix} 1 & 0 & 0 & 0 & 0 \\ 0 & 1 & 0 & 0 & 0 \\ 0 & 0 & 2 & 0 & 0 \\ 0 & 0 & 0 & 6 & 0 \\ 0 & 0 & 0 & 0 & 24 \end{pmatrix}$$

```
In[4] :=
```



- Define multivariate generating function

$$\omega(\mathbf{x}) = \frac{1}{(2\pi)^{d/2}} \exp\left(-\frac{\mathbf{x}^2}{2}\right), \mathbf{H}^{(n)}(\mathbf{x}) = (-1)^n \frac{1}{\omega(\mathbf{x})} \nabla^{(n)} \omega(\mathbf{x}), \nabla^{(n)} = \nabla \otimes \dots \otimes \nabla$$

$$\nabla^{(1)} = \nabla = \begin{pmatrix} \frac{\partial}{\partial x_1} \\ \vdots \\ \frac{\partial}{\partial x_d} \end{pmatrix}, \nabla^{(2)} = \nabla \otimes \nabla = \begin{pmatrix} \frac{\partial^2}{\partial x_1^2} & \cdots & \frac{\partial^2}{\partial x_1 \partial x_d} \\ \vdots & \ddots & \vdots \\ \frac{\partial^2}{\partial x_d \partial x_1} & \cdots & \frac{\partial^2}{\partial x_d^2} \end{pmatrix}$$

```
In[4] := omega[x_,y_,z_] = Exp[-(x^2+y^2+z^2)/2]/(2Pi)^(3/2)
```

$$\frac{e^{\frac{1}{2}(-x^2-y^2-z^2)}}{2\sqrt{2}\pi^{3/2}}$$

```
In[5] :=
```



```
In[5] := H[n_,i_,j_,k_] :=
  Simplify[(-1)^n/omega[x,y,z] D[omega[x,y,z],{x,i},{y,j},{z,k}]];
H[n_] := H[n] = DeleteCases[Flatten[Table[If[i+j+k==n,H[n,i,j,k],{}],
  {k,0,n},{j,0,n},{i,0,n}]],{}];
H[0]
```

{1}

```
In[6] := H[1]
```

{x, y, z}

```
In[7] := H[2]
```

{x² - 1, x y, y² - 1, x z, y z, z² - 1}

```
In[8] := H[3]
```

{x(x² - 3), (x² - 1)y, x(y² - 1), y(y² - 3), (x² - 1)z, x y z, (y² - 1)z, x(z² - 1), y(z² - 1), z(z² - 3)}

```
In[9] :=
```

$$\int \omega(\mathbf{x}) H_{\alpha}^{(m)}(\mathbf{x}) H_{\beta}^{(n)}(\mathbf{x}) d\mathbf{x} = n_x! n_y! n_z! \delta_{mn}^{(2)} \delta_{\alpha\beta}^{(m+n)}, \delta_{\alpha\beta}^{(m+n)} = 1 \text{ if } \alpha \text{ is permutation of } \beta$$

$$\alpha = (x, x, y) \Rightarrow (n_x, n_y, n_z) = (2, 1, 0).$$



- Seek Hermite expansion of equilibrium function ($d = 3$)

$$f^{\text{eq}}(\boldsymbol{\xi}; \rho, \mathbf{u}, \theta) = \frac{\rho}{(2\pi\theta)^{d/2}} \exp\left[-\frac{(\boldsymbol{\xi} - \mathbf{u})^2}{2\theta}\right] = \omega(\boldsymbol{\xi}) \sum_{n=0}^{\infty} \frac{1}{n!} \mathbf{a}^{(n)}(\rho, \mathbf{u}, \theta) \cdot \mathbf{H}^{(n)}(\boldsymbol{\xi})$$

$$\mathbf{a}^{(n)}(\rho, \mathbf{u}, \theta) = \int f^{\text{eq}}(\boldsymbol{\xi}; \rho, \mathbf{u}, \theta) \mathbf{H}^{(n)}(\boldsymbol{\xi}) d\boldsymbol{\xi}$$

- $f^{\text{eq}}(\boldsymbol{\xi})$ of same form as $\omega(\boldsymbol{\xi}) = \exp(-\boldsymbol{x}^2/2) / (2\pi)^{d/2}$

$$f^{\text{eq}}(\boldsymbol{\xi}; \rho, \mathbf{u}, \theta) = \frac{\rho}{(2\pi\theta)^{d/2}} \exp\left[-\frac{(\boldsymbol{\xi} - \mathbf{u})^2}{2\theta}\right] = \frac{\rho}{\theta^{d/2}} \omega\left(\frac{\boldsymbol{\xi} - \mathbf{u}}{\sqrt{\theta}}\right)$$

- Compute coefficients $\mathbf{a}^{(n)}(\rho, \mathbf{u}, \theta)$, $\boldsymbol{\eta} = (\boldsymbol{\xi} - \mathbf{u}) / \sqrt{\theta}$ (scale to thermal velocity)

$$\mathbf{a}^{(n), \text{eq}}(\rho, \mathbf{u}, \theta) = \frac{\rho}{\theta^{d/2}} \int \omega\left(\frac{\boldsymbol{\xi} - \mathbf{u}}{\sqrt{\theta}}\right) \mathbf{H}^{(n)}(\boldsymbol{\xi}) d\boldsymbol{\xi} = \rho \int \omega(\boldsymbol{\eta}) \mathbf{H}^{(n)}(\sqrt{\theta} \boldsymbol{\eta} + \mathbf{u}) d\boldsymbol{\eta}$$



$$\mathbf{a}^{(n),\text{eq}}(\rho, \mathbf{u}, \theta) = \frac{\rho}{\theta^{d/2}} \int \omega\left(\frac{\boldsymbol{\xi} - \mathbf{u}}{\sqrt{\theta}}\right) \mathbf{H}^{(n)}(\boldsymbol{\xi}) d\boldsymbol{\xi} = \rho \int \omega(\boldsymbol{\eta}) \mathbf{H}^{(n)}(\sqrt{\theta} \boldsymbol{\eta} + \mathbf{u}) d\boldsymbol{\eta}$$

```
In[9] := aeq[n_] := aeq[n] = Module[{Hn, an, nH, HnE, inf, st},
  Hn=H[n]; nH=Length[Hn]; inf=Infinity; st=Sqrt[theta];
  HnE = Hn /. {x->st r + u, y->st s + v, z->st t + w};
  an=Table[Integrate[ omega[r,s,t] HnE[[k]],
    {r,-inf,inf},{s,-inf,inf},{t,-inf,inf}],{k,1,nH}];
  Return[rho an] ];
Aeq={aeq[0],aeq[1]}
```

```
{{rho}, {rho u, rho v, rho w}}
```

```
In[10] := AppendTo[Aeq, aeq[2]]
```

```
{{rho}, {rho u, rho v, rho w}, {rho (theta + u^2 - 1), rho u v, rho (theta + v^2 - 1), rho u w, rho v w, rho (theta + w^2 - 1)}}
```

```
In[11] := Simplify[Aeq[[3]][[1]]+Aeq[[3]][[3]]+Aeq[[3]][[6]]]
```

```
rho (3 theta + u^2 + v^2 + w^2 - 3)
```

```
In[12] :=
```



```
In[12] := AppendTo[Aeq, aeq[3]] ;  
Aeq[[1]]
```

$$\{\rho\}$$

```
In[13] := Aeq[[2]]
```

$$\{\rho u, \rho v, \rho w\}$$

```
In[14] := Aeq[[3]]
```

$$\{\rho(\theta + u^2 - 1), \rho u v, \rho(\theta + v^2 - 1), \rho u w, \rho v w, \rho(\theta + w^2 - 1)\}$$

```
In[15] := Aeq[[4]]
```

$$\{\rho u(3\theta + u^2 - 3), \rho v(\theta + u^2 - 1), \rho u(\theta + v^2 - 1), \rho v(3\theta + v^2 - 3), \rho w(\theta + u^2 - 1), \rho u v w, \rho w(\theta + v^2 - 1), \rho u(\theta + w^2 - 1), \rho v(\theta + w^2 - 1), \rho w(3\theta + w^2 - 3)\}$$

```
In[16] :=
```

First three terms of Hermite expansion correspond to conserved quantities



```
In[16] := vRules = {x->xi, y->eta, z->zeta};
          feqH[xi_,eta_,zeta_] = omega[xi,eta,zeta] Sum[aeq[n].H[n]/n! /. vRules, {n,0,2}]
```

$$e^{\frac{1}{2}(-\eta^2 - \xi^2 - \zeta^2)} \left(\frac{1}{2} ((\eta^2 - 1) \rho (\theta + v^2 - 1) + \eta \rho u v \xi + \eta \rho v w \zeta + \rho (\xi^2 - 1) (\theta + u^2 - 1) + \rho (\zeta^2 - 1) (\theta + w^2 - 1) + \rho u w \xi \zeta) + \eta \rho v + \rho u \xi + \rho w \zeta + \rho \right) / (2 \sqrt{2} \pi^{3/2})$$

```
In[18] := inf=Infinity;
          Integrate[feqH[xi,eta,zeta],{xi,-inf,inf},{eta,-inf,inf},{zeta,-inf,inf}]
```

ρ

```
In[19] := Integrate[feqH[xi,eta,zeta] xi,{xi,-inf,inf},{eta,-inf,inf},{zeta,-inf,inf}]
```

ρu

```
In[20] := Integrate[feqH[xi,eta,zeta] eta,{xi,-inf,inf},{eta,-inf,inf},{zeta,-inf,inf}]
```

ρv

```
In[21] := Integrate[feqH[xi,eta,zeta] zeta,{xi,-inf,inf},{eta,-inf,inf},{zeta,-inf,inf}]
```

ρw

```
In[22] :=
```



```
In[23]:= Integrate[feqH[xi,eta,zeta] xi xi,{xi,-inf,inf},{eta,-inf,inf},{zeta,-inf,inf}]
```

$$\rho(\theta + u^2)$$

```
In[24]:= Integrate[feqH[xi,eta,zeta] eta eta,{xi,-inf,inf},{eta,-inf,inf},{zeta,-inf,inf}]
```

$$\rho(\theta + v^2)$$

```
In[25]:= Integrate[feqH[xi,eta,zeta] zeta zeta,{xi,-inf,inf},{eta,-inf,inf},{zeta,-inf,inf}]
```

$$\rho(\theta + w^2)$$



```
In[26] := Integrate[feqH[xi,eta,zeta] xi eta,{xi,-inf,inf},{eta,-inf,inf},{zeta,-inf,inf}]
```

$$\frac{\rho u v}{2}$$

```
In[27] := Integrate[feqH[xi,eta,zeta] xi zeta,{xi,-inf,inf},{eta,-inf,inf},{zeta,-inf,inf}]
```

$$\frac{\rho u w}{2}$$

```
In[28] := Integrate[feqH[xi,eta,zeta] eta zeta,{xi,-inf,inf},{eta,-inf,inf},{zeta,-inf,inf}]
```

$$\frac{\rho v w}{2}$$



- Verify $\mathbf{a}^{(n),\text{eq}} = \mathbf{a}^{(n),\text{eq}} \otimes \mathbf{u}$, $a^{(0),\text{eq}} = \rho$

```
In[29] := V={u,v,w}; a0[1] = aeq[1] /. rho -> Subscript[rho,0]
```

```
{rho u, rho v, rho w}
```

```
In[30] := a0[2] = Outer[Times,a0[1],V]
```

```

$$\begin{pmatrix} \rho_0 u^2 & \rho_0 u v & \rho_0 u w \\ \rho_0 u v & \rho_0 v^2 & \rho_0 v w \\ \rho_0 u w & \rho_0 v w & \rho_0 w^2 \end{pmatrix}$$

```

```
In[31] := aeq[2] /. theta->1
```

```
{rho u^2, rho u v, rho v^2, rho u w, rho v w, rho w^2}
```

```
In[32] :=
```



- Third-order

```
In[32] := a0[3] = DeleteDuplicates[Flatten[Outer[Times, a0[2], V]]]
```

```
{ρ0 u3, ρ0 u2 v, ρ0 u2 w, ρ0 u v2, ρ0 u v w, ρ0 u w2, ρ0 v3, ρ0 v2 w, ρ0 v w2, ρ0 w3}
```

```
In[33] := aeq[3] /. theta->1
```

```
{ρ u3, ρ u2 v, ρ u v2, ρ v3, ρ u2 w, ρ u v w, ρ v2 w, ρ u w2, ρ v w2, ρ w3}
```

```
In[34] :=
```



- $\mathbf{a}^{(n),\text{eq}}(\rho, \mathbf{u}, \theta) = \int f^{\text{eq}}(\boldsymbol{\xi}) \mathbf{H}^{(n)}(\boldsymbol{\xi}) d\boldsymbol{\xi}$, $f^{\text{eq},N}(\boldsymbol{\xi}) = \omega(\boldsymbol{\xi}) \rho Q(\mathbf{u}, \theta, \boldsymbol{\xi}) \Rightarrow$

$$\mathbf{a}^{(n),\text{eq}}(\rho, \mathbf{u}, \theta) = \rho \int \omega(\boldsymbol{\xi}) Q(\mathbf{u}, \theta, \boldsymbol{\xi}) \mathbf{H}^{(n)}(\boldsymbol{\xi}) d\boldsymbol{\xi} = \sum_{i=0}^{[(n+2)/2]} w_i Q(\mathbf{u}, \theta, \boldsymbol{\xi}_i) \mathbf{H}^{(n)}(\boldsymbol{\xi}_i)$$

In[34] := H[2, x]

$$x^2 - 1$$

In[35] :=

$$\int_{-\infty}^{\infty} w(x) f(x) dx \cong \sum_{i=0}^{q-1} w_i f(x_i), \quad \int_{-\infty}^{\infty} w(x) p_{2n-1}(x) dx = \sum_{i=0}^{n-1} w_i p_n(x_i)$$



- Consider a discrete velocity set $\{\xi_0, \dots, \xi_q\}$. The Hermite expansion of the Maxwell-Boltzmann distribution up to second order $N = 2$ is

$$f^{\text{eq}}(\xi; \rho, \mathbf{u}, \theta) \cong f^{\text{eq}, N}(\xi; \rho, \mathbf{u}, \theta) = \omega(\xi) \rho Q(\xi; \mathbf{u}, \theta)$$

- Evaluate the $N = 2$ approximation for discrete velocity ξ_i

$$f_i^{\text{eq}} = \rho w_i Q(\xi_i; \mathbf{u}, \theta), w_i \equiv \omega(\xi_i)$$

```
In[39] := Q[xi_, eta_, zeta_] = Simplify[feqH[xi, eta, zeta] / omega[xi, eta, zeta] / rho]
```

$$\frac{1}{2} (\eta^2 (\theta + v^2 - 1) + \eta v (u \xi + w \zeta + 2) + \theta (\xi^2 + \zeta^2 - 3) + u^2 \xi^2 - u^2 + u w \xi \zeta + 2 u \xi - v^2 + w^2 \zeta^2 - w^2 + 2 w \zeta - \xi^2 - \zeta^2 + 5)$$

```
In[47] := Q[xi, 0, 0] /. theta -> 1
```

$$\frac{1}{2} (u^2 \xi^2 - u^2 + 2 u \xi - v^2 - w^2 + 2)$$

```
In[48] :=
```



- For a discrete set of velocities $\{\mathbf{c}_i\}_{i=0,1,\dots,Q}$ the discrete Boltzmann equation is

$$\partial_t f_i + c_{i\alpha} \partial_\alpha f_i = \Omega(f_i), i = 0, 1, \dots, Q,$$

$$\partial_t f_i + (\mathbf{c}_i \cdot \nabla) f_i = \Omega(f_i)$$

- Step 1 of operator splitting, $\partial_t f_i + (\mathbf{c}_i \cdot \nabla) f_i = 0,$

$$f_i(t + \delta t, \mathbf{x}) = f_i(t, \mathbf{x} - \mathbf{c}_i \delta t)$$

- Step 2 of operator splitting

$$\partial_t f_i = \Omega(f_i)$$



$$\frac{df_i}{dt} = -\frac{(f_i - f_i^{(0)})}{\tau} \Rightarrow f_i(1) = e^{-1/\tau} f_i(0) + f_i^{\text{eq}}(1 - e^{-1/\tau}) = e^{-1/\tau}(f_i(0) - f_i^{\text{eq}}) + f_i^{\text{eq}}$$

$$f_i(1) = \left(1 - \frac{1}{\tau} + \frac{1}{2!\tau^2} - \frac{1}{3!\tau^3} + \dots\right)(f_i(0) - f_i^{\text{eq}}) + f_i^{\text{eq}} \cong \left(1 - \frac{1}{\tau}\right)(f_i(0) - f_i^{\text{eq}}) + f_i^{\text{eq}}$$

$$f_i(1) = f_i(0) + \frac{1}{\tau}(f_i^{\text{eq}} - f_i(0))$$

In [50] := Simplify[DSolve[{fi'[t]==-(fi[t]-fieg)/tau,fi[0]==fi0},fi[t],t]]

$$\left\{\left\{f_i(t) \rightarrow e^{-\frac{t}{\tau}}(f_i(0) + f_i^{\text{eq}}(e^{t/\tau} - 1))\right\}\right\}$$

In [51] :=

$$y' = f(y) = \frac{\partial f}{\partial y} y, |1 + z| \leq 1, z = \lambda \delta t, y \sim e^{\lambda t}$$